

NMR Spectroscopy

Nuclear magnetic Resonance is a branch of spectroscopy in which radio frequency waves induce transitions between magnetic energy levels of nuclei of a molecule. The magnetic energy levels are created by keeping the nuclei in a magnetic field.

Without the magnetic field the spin states of nuclei are degenerate i.e. possess the same energy, and energy level transition is not possible. When a magnetic field is applied, the separate levels and radio frequency radiation can cause transitions between these energy levels.

(NMR spectroscopy is most often concerned with nuclei with $I = 1/2$) Examples of such nuclei are ^1H , ~~^{31}P~~ ^{31}P and ^{19}F . Spectra cannot be obtained from nuclei with $I = 0$. In special cases, spectra can be obtained from nuclei when $I \geq 1$.

(Nuclear magnetic Resonance is a powerful tool for investigating nuclear structure.)

In nuclear magnetic resonance spectroscopy radio frequency waves induce transitions between magnetic energy levels of nuclei of a molecule.

The magnetic energy levels are created by keeping the nuclei in a magnetic field.

Spin and an applied field

We have seen that all electrons have a spin of $\frac{1}{2}$ they have an angular momentum

$$\text{of } \sqrt{\frac{1}{2}(\frac{1}{2}+1)} \frac{h}{2\pi} = \frac{\sqrt{3}}{2} h$$

Many nuclei also possess spin

Let us consider the simplest nucleus

hydrogen atom which consists of only one proton

The proton has a spin of $\frac{1}{2}$

Thus if a particular nucleus is composed of p protons and n neutrons its total mass is $p+n$

its total charge is p and the total

spin will be a vector combination of

$p+n$ spins each of magnitude $\frac{1}{2}$

Then the spin of hydrogen nucleus is $\frac{1}{2}$ since it contains one proton only, deuterium an isotope of hydrogen containing one proton and one neutron (${}^2\text{H}$) might have a spin of 1 or 0 depending on whether the proton and neutron spins are parallel or opposed.

1. Nuclei with both p and n even have zero spin (${}^4\text{He}$, ${}^{12}\text{C}$, ${}^{16}\text{O}$ etc)
2. Nuclei with both p and n odd ($p+n = \text{even}$) have integral spin (${}^2\text{H}$, ${}^{14}\text{N}$, ${}^{10}\text{B}$)
3. Nuclei with odd mass ~~no~~ have half-integral spins (${}^1\text{H}$, ${}^{15}\text{N}$)

The spin of a nucleus is usually given by the symbol I , called the spin or angular momentum of the nucleus

$$I = \sqrt{I(I+1)} \frac{h}{2\pi} = \sqrt{I(I+1)} \text{ units.}$$

The angular momentum vector cannot point in any arbitrary direction, but can point only along a particular reference direction so that the

The components have half integral or integral value along that direction

$$I_2 = (1, 1, -1, \dots, 0, 1, -(I-1) - 1) \quad \left[\begin{array}{l} \text{for} \\ I \text{ integral} \end{array} \right]$$

or

$$I_2 = (1, 1-1, \dots, 1/2, -1/2, \dots, -1)$$

for I half integral

Quantum Mechanical description of NMR

If a nucleus has spin, it behaves as a spinning finite ~~partial~~ distribution of charge.

One of the most important properties of a nucleus is its spin I or intrinsic spin angular momentum $I\hbar$. This gives rise to a magnetic moment μ to the nucleus.

The magnetic moment μ associated with the spin angular momentum is given by

$$\mu = \gamma I \hbar \quad \text{--- (1)}$$

Here γ is a scalar called the gyromagnetic ratio.

and it may take +ve or -ve value.

An alternate expression for the magnetic moment is

$$\mu = g_N \mu_N I \quad \text{--- (2)}$$

where g_N - nuclear g factor. It is a dimensionless quantity

μ_N is the nuclear magneton

$$\mu_N = \frac{e \hbar}{2 m_p} \quad \text{--- (3)}$$

Resonance condition:

When a nucleus of magnetic moment μ is placed in a magnetic field B_0 ,

The interaction energy

$$E = -\mu \cdot B_0$$
$$= -\mu B_0 \cos\theta$$

Eqn $E = \frac{-\mu B_0 m_I}{I}$ ——— 4

m_I - Projection of spin vector on magnetic field direction

m_I can have

m_I - Projection of spin vector I in any direction

m_I can have $(2I+1)$ values, so there are $2I+1$ equally spaced energy levels.

The energy separation between any two adjacent level is given by

$$|\Delta E| = \mu B_0 \frac{\Delta m_I}{I} = \frac{\mu B_0}{I}$$
$$= g_N \mu_N B_0$$

The basis of NMR experiment is to introduce a transition from a lower level to the next higher level. If ν is the frequency of the electromagnetic radiation that induces transitions between adjacent levels, Bohr's condition is then

$$h\nu = \frac{\mu B_0}{I} = g_N \mu_N B_0$$

$$\nu = \frac{g_N \mu_N B_0}{h} \quad \text{--- (5)}$$

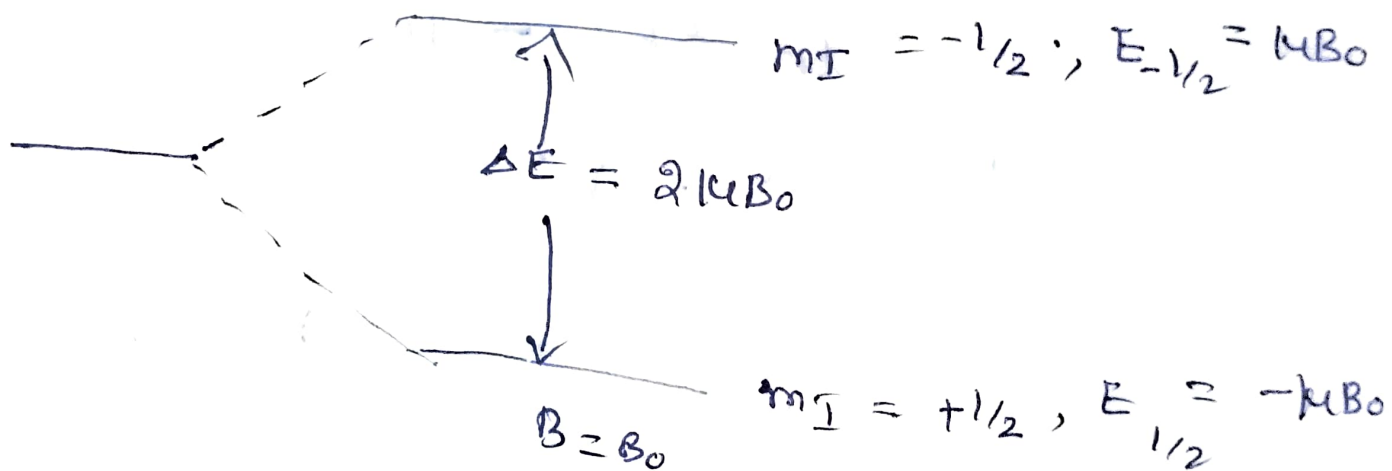
Eqn (5) is the resonance condition.

For a spin $1/2$ system we will have 2 states, one corresponding to $m_I = 1/2$ and the other for $m_I = -1/2$

These energies are $E_{1/2} = -\mu B_0$ and $E_{-1/2} = +\mu B_0$ — (6)

The resonance condition reduces to

$$h\nu = 2\mu B_0$$



Coupling constants

The distance between the peaks in a given multiplet is a measure of the magnitude of splitting effect. It is referred to as coupling constant and is denoted by the symbol J . The numerical value of J is expressed in Hz or cps.

Unlike the chemical shifts, the values of J are independent of the applied field strength and depend only on the molecular structure.