$$\lambda_{(mA)} = \frac{1}{v(cm^{-1})} = 2.341 \times 10^{*}$$

Pure Rotational Raman Spectra. The energy of a rigid rotator will be given by

$$E_{r} = \frac{h^{2}}{8\pi^{2}l}J(J+1)$$
(4.17)

The selection rule for rotational Raman spectrum is as follows :

hes $\Delta J=0$, the scattered Raman radiation will be of the same frequency as that of incident light (Rayseign's intring). The transition $\Delta J=\pm 2$ gives Stoke's lines (longer wavelength) whereas $\Delta J = -2$ gives the antitake's lines (shorter wavelength).

When AJ=+2, the values of rotational Raman shifts (Stoke's lines) will be given by

$$\Delta \bar{v} = \frac{h}{8\pi^2 lc} [(J+2) (J+3) - J(J+1)]$$

$$= 2B(2J+3) \text{ where } B = \frac{h}{8\pi^2 lc}$$
(4.19)

then AJ = -2, the values of rotational Raman shifts (anti-Stoke's lines) will be given by

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 $\Delta \tilde{v} = -2B (2J+3)$ On combining equations (4.19) and (4.20), the Raman shift can be put in the form of the for

$$\Lambda \bar{v} = \pm 2B(2J + 3)$$
 where $J = 0, 1, 2, 3...$

The wave numbers of the corresponding spectral lines will be given by

$$v = v_{ex} + \Delta v$$

Where v_{ex} is the wave number of exciting radiation,

The transitions and the Raman spectrum are shown in Fig. 4.5.



Fig. 4.5 : Rotational Raman Spectrum of a Rigid Diatomic Molecule.

From the Fig. 4.5 it can be seen that the frequency separation of successive lines is 4B cm⁻¹ when it is 2B cm⁻¹ in the far-infrared spectra while on substituting J=0 in equation (4.21) we observe that separation of the first line from the exciting line will be 6B cm⁻¹.

4. Vibrational-Rotational Raman Spectra. It is possible to predict theoretically that vibrational a rotational transitions may occur simultaneously in a Raman transition. In such a case, the size rules governing such simultaneous changes in rotational and vibrational energy will be

 $\Delta J = 0, \pm 2$ and $\Delta v = \pm 1$

For a diatomic molecule, the vibration-rotational energy levels are given by

$$\mathbf{E}_{rv} = hc \left[\overline{w_e} \left(v + \frac{1}{2} \right) - \overline{w_e} x \left(v + \frac{1}{2} \right)^2 \right] + Bhc \ \mathbf{J} (\mathbf{J} + 1)$$

where v = 0, 1, 2,

and J = 0, 1, 2,.....

Or, in terms of frequency,

$$\overline{v} = \overline{w_e} \left(v + \frac{1}{2} \right) - \overline{w_e} x \left(v + \frac{1}{2} \right)^2 + B J (J+1)$$

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Applying selection rules for Raman effect, i.e., J-0

$$\Delta J = 0; \Delta \overline{v} (Q) = \overline{v_0} (\text{for all } J)$$

$$\Delta J = +2; \Delta \overline{v} (S) = \overline{v_0} + B(4J + 6)(\text{for } J = 0, 1, 2, ...)$$

$$\Delta J = -2; \Delta \overline{v} (O) = \overline{v_0} - B(4J + 6)(\text{for } J = 2, 3, 4, ...)$$

where $\overline{v_0}$ is written for $\overline{w_e}$ (1-2x), O, Q and S refer to the O branch lines, Q branch lines, and S branch lines, respectively. The corresponding Stoke's lines will occur at wave numbers represented by

(i)
$$\overline{v}(Q) = \overline{v_{ex}} - \overline{v}(Q) = \overline{v_{ex}} - \overline{v_0}$$
 (for all values of J)
(ii) $\overline{v}(S) = \overline{v_{ex}} - \overline{v}(S) = \overline{v_{ex}} - \overline{v_0} - B$ (4J + 6), where J = 0,1,2,...
(jii) $\overline{v}(O) = \overline{v_{ex}} - \overline{v}(O) = \overline{v_{ex}} - \overline{v_0} + B$ (4J + 6), where J = 2,3,4,...

The various Raman rotation-vibrations transitions are shown in Fig. 4.5. In the Raman spectrum Q branch will be present whereas it is absent in infrared spectrum. Both S and O branches are somewhat smilar to the R and P branches in infrared spectroscopy.



Let us now briefly consider the rotational transitions accompanying a vibrational Raman transition of a diatomic molecule. The selection rules now involve the changes in both the vibrational and the rotational quantum numbers. For a diatomic molecule, the selection rules are

$$\Delta v = \pm 1 \tag{87}$$

Since at room temperature most of the molecules are in the ground vibrational state (v=0), only the vibrational transition, v=0 to v=1, is of interest. The transitions which $\Delta J=0$ form a *Q*-branch; those with $\Delta J=+2$ form an *S*-branch and those with $\Delta J=-2$ form an *O*-branch. The rotational Raman transitions accompanying a $0\rightarrow 1$ vibrational transition are shown in Fig. 4.6. Here, $\Delta \overline{v}$ measures the displacement from the exciting mercury line. The *Q*-branch exhibits an intense 'narrow line' at \overline{v}_{ex} that is usually unresolved while the *S* and *O* branches form weak wings which extend to lower and higher were numbers, respectively, from the intense narrow line.

2.91